

THEORY ISSUES IN $e^+e^- \rightarrow H$ MEASUREMENTS

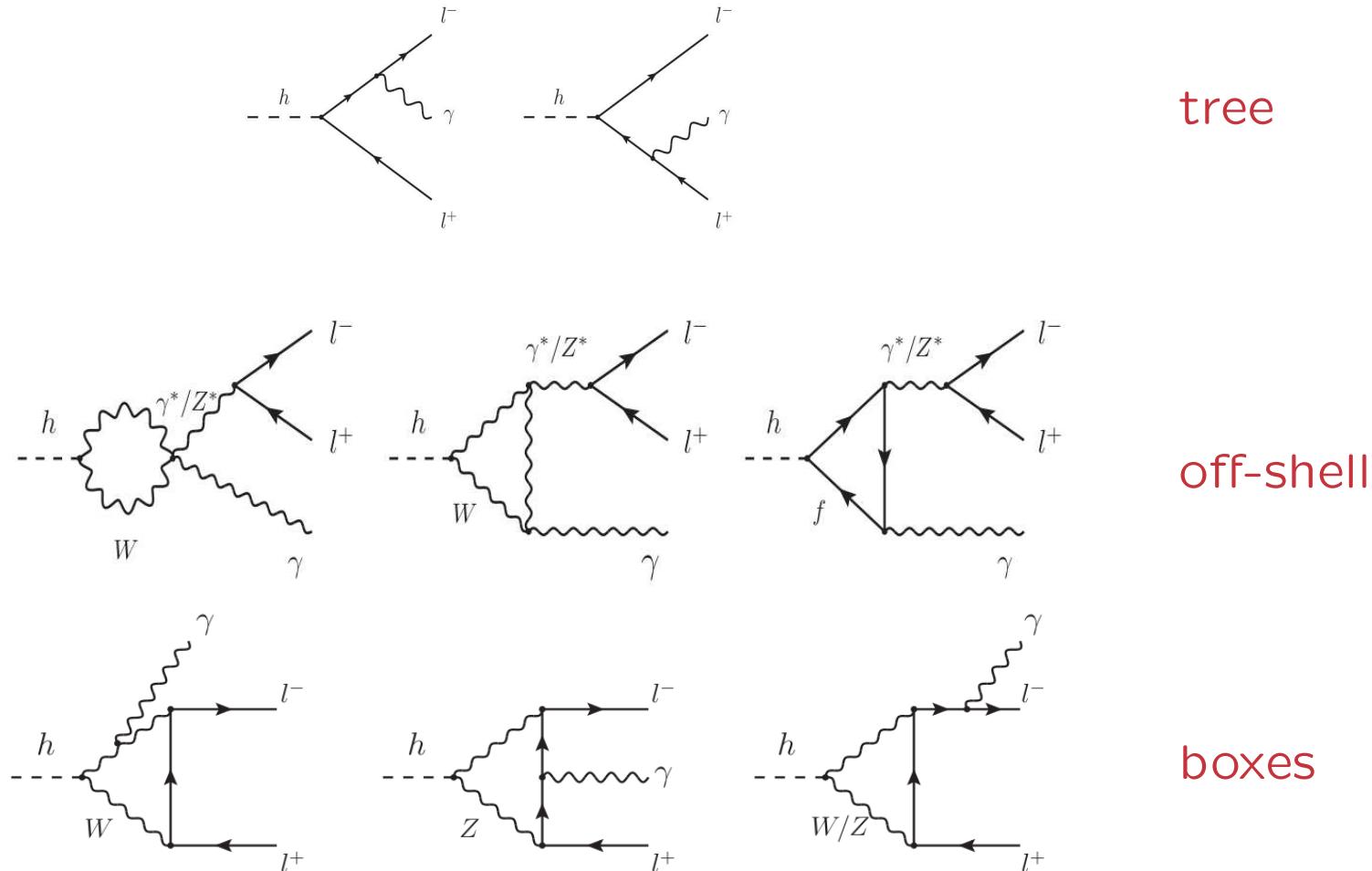
Michael Spira (PSI)

I Dalitz Decays

II Higgs production

III Conclusions

I HIGGS DALITZ DECAYS

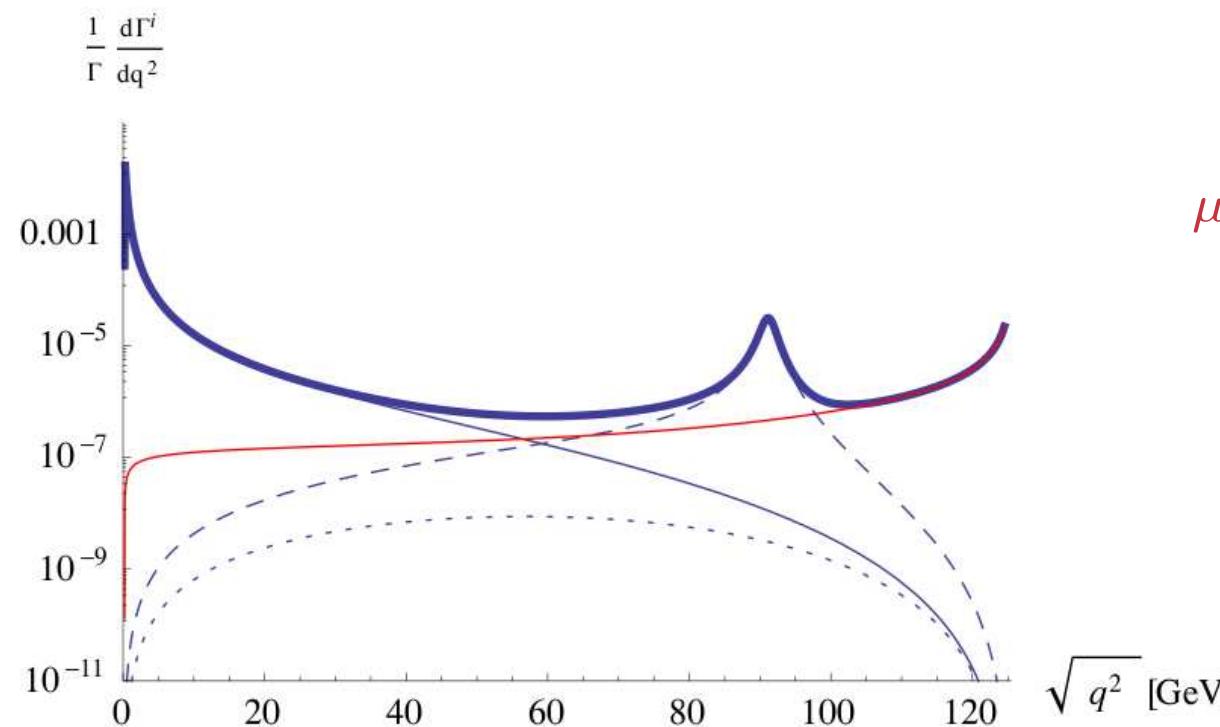
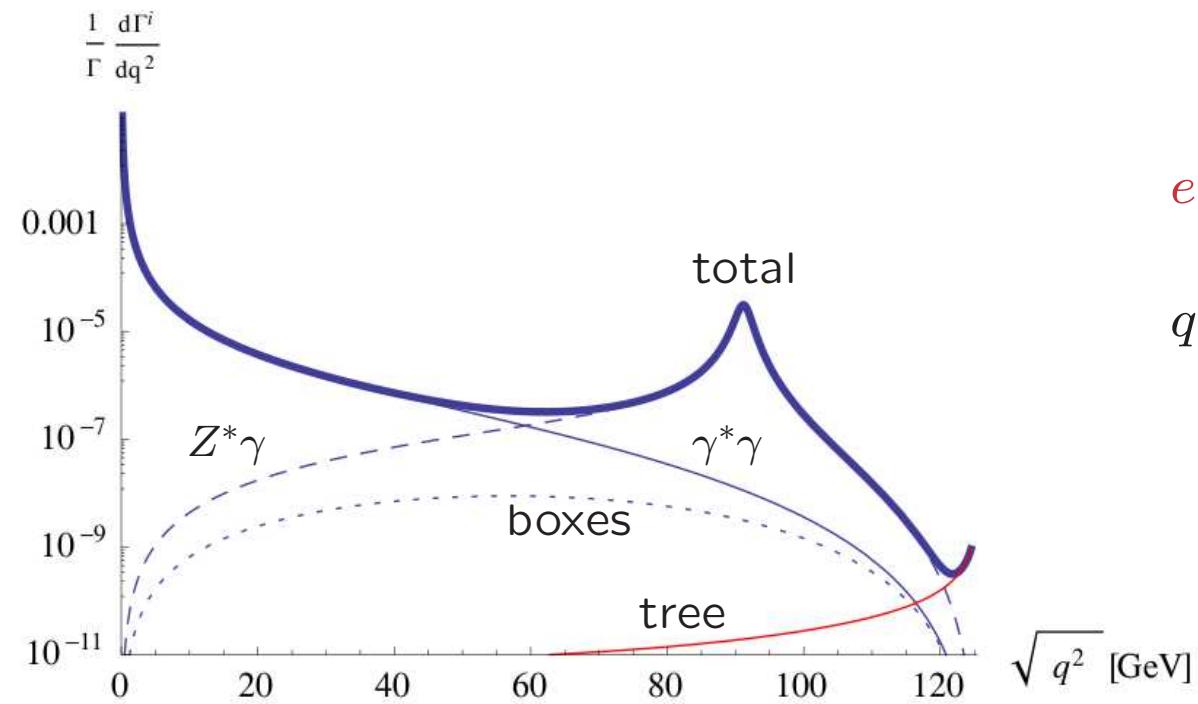


Abbasabadi, Bowser-Chao, Dicus, Repko
Sun, Chang, Gao
Passarino
Kachanovich, Nierste, Nišandžić

$$\frac{\Gamma(h\rightarrow \gamma e^+e^-)}{\Gamma(h\rightarrow \gamma\gamma)}=5.7\%$$

$$\frac{\Gamma(h\rightarrow \gamma \mu^+\mu^-)}{\Gamma(h\rightarrow \gamma\gamma)}=5.8\%~~~~~(E_\gamma>1~{\rm GeV})$$

$$\frac{\Gamma(h\rightarrow \gamma \tau^+\tau^-)}{\Gamma(h\rightarrow \gamma\gamma)}=3.04$$



Higgs Boson Decays: Theoretical Status

Michael Spira (PSI, Villigen). May 13, 2019. 12 pp.

Published in CERN Yellow Reports: Monographs 3 (2020) 123-134

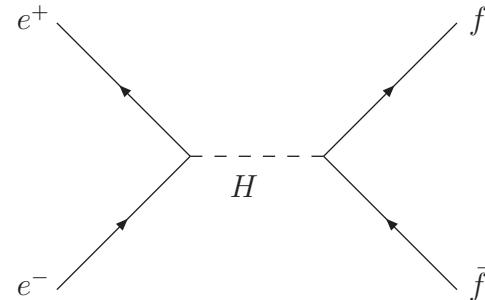
DOI: 10.23731/CYRM-2020-003.123

[†]It should be noted that the end point in the $e^+e^-\gamma$ case is 4–5 orders of magnitude smaller than the photon and Z-exchange contributions thus making it impossible to determine the electron Yukawa coupling. The same conclusion is also valid for the reverse process $e^+e^- \rightarrow H\gamma$ so that the *s*-channel line-shape measurement proposed in Ref. [89] will not be sensitive to the electron Yukawa coupling but dominated by the loop-induced contribution with an additional photon.

- [89] S. Jadach, R. A. Kycia, Lineshape of the Higgs boson in future lepton colliders, Phys. Lett. B755 (2016) 58–63. [arXiv:1509.02406](https://arxiv.org/abs/1509.02406), doi:[10.1016/j.physletb.2016.01.065](https://doi.org/10.1016/j.physletb.2016.01.065).

II HIGGS PRODUCTION

$$e^+ e^- \rightarrow H \rightarrow f\bar{f}$$

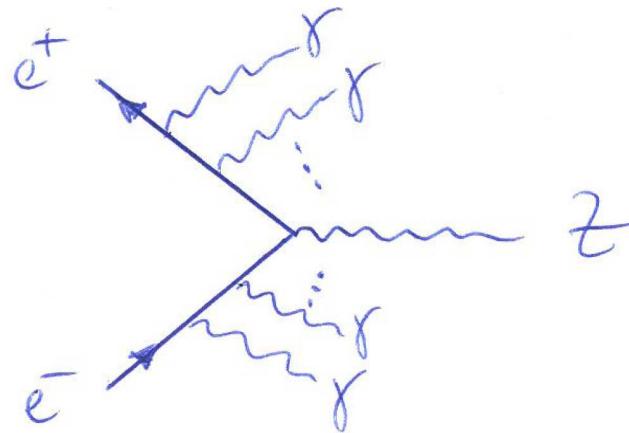


$$\sigma_0(e^+ e^- \rightarrow H \rightarrow f\bar{f}) = N_{cf} \frac{m_e^2 m_f^2}{16\pi v^4} \frac{s}{(s - M_H^2)^2 + M_H^2 \Gamma_H^2}$$

$$\Gamma(H \rightarrow f\bar{f}) = N_{cf} \frac{m_f^2}{8\pi v^2} M_H$$

$$\Rightarrow \sigma_0(e^+ e^- \rightarrow H) = \frac{m_e^2 \Gamma_H}{2v^2 M_H} \frac{s}{(s - M_H^2)^2 + M_H^2 \Gamma_H^2}$$

initial state radiation (ISR):

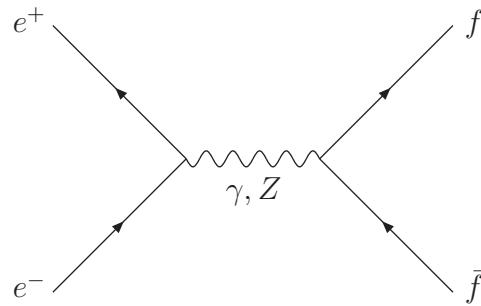


$$\sigma(e^+ e^- \rightarrow H) = \int_0^1 d\nu \rho(\nu) \sigma_0[\hat{s} = (1-\nu)s]$$

↑
YFS resummation kernel

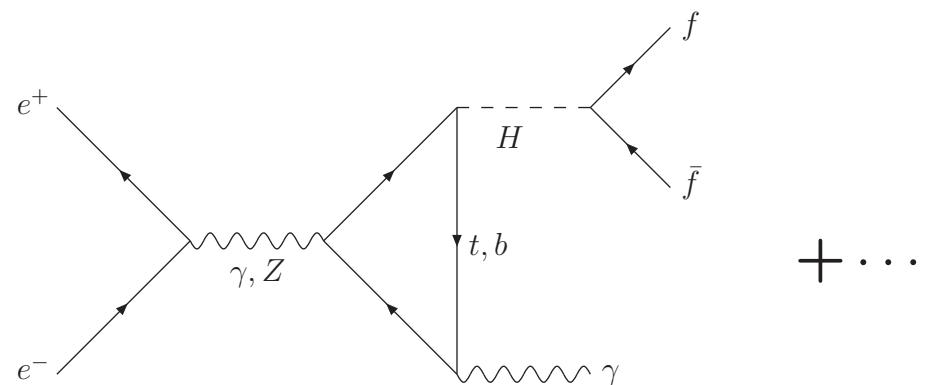
Jadach, Kycia

$$e^+ e^- \rightarrow f\bar{f}$$



$$\sigma(e^+ e^- \rightarrow f\bar{f}) \sim (10^3 \dots 10^4) fb$$

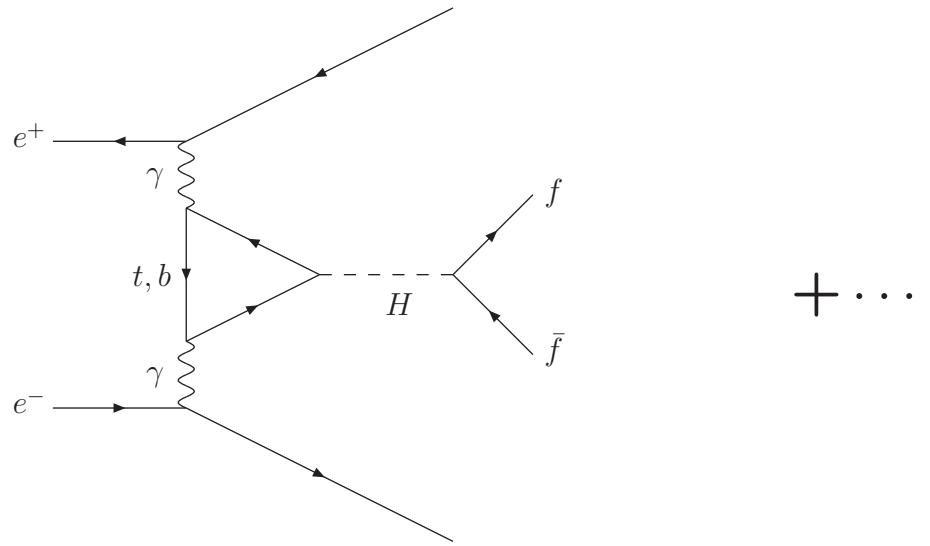
inverse Dalitz: $e^+ e^- \rightarrow H\gamma \rightarrow f\bar{f}\gamma$



$$\sigma(e^+ e^- \rightarrow H\gamma) = \frac{2\pi\Gamma_H}{sM_H} \sum |\mathcal{M}|^2 \frac{p^2}{(p^2 - M_H^2)^2 + M_H^2\Gamma_H^2} dPS_3$$

Abbasabadi, Bowser-Chao, Dicus, Repko

photon fusion: $\gamma\gamma \rightarrow H \rightarrow f\bar{f}$



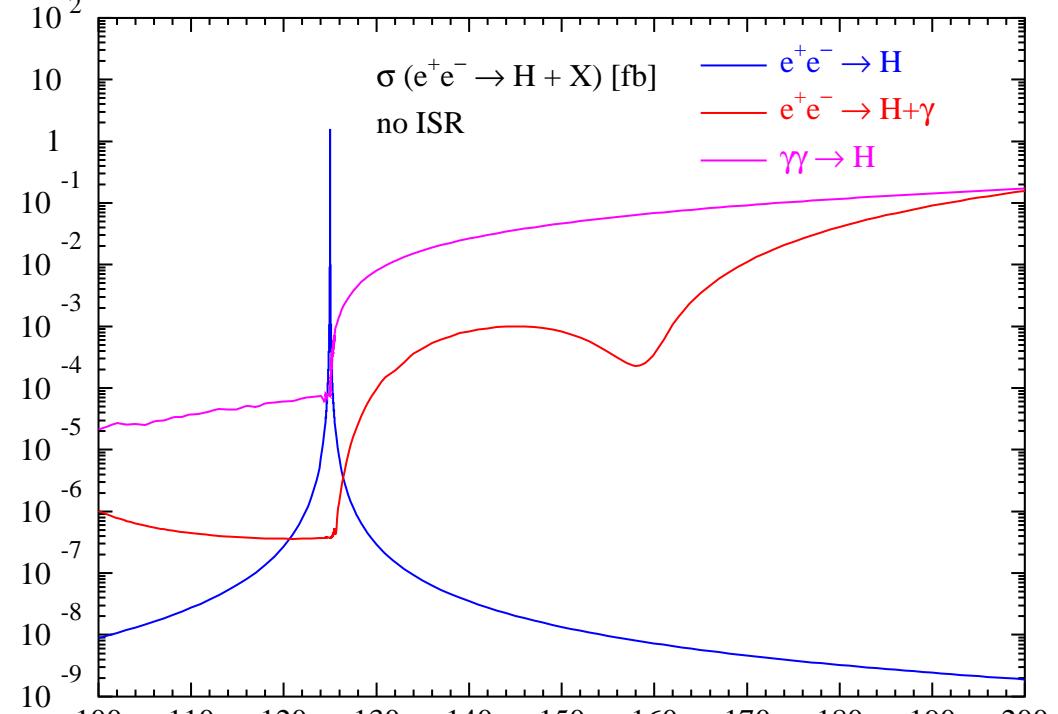
$$\sigma(e^+e^- \rightarrow He^+e^-) \approx \int_{\tau_0}^1 d\tau \int_{\tau}^1 \frac{dx}{x} f_{\gamma/e}(x) f_{\gamma/e}\left(\frac{\tau}{x}\right) \sigma_{\gamma\gamma \rightarrow H}(\hat{s} = \tau s)$$

$$\sigma_{\gamma\gamma \rightarrow H} = \frac{\Gamma_H}{4\hat{s}M_H} \sum |\mathcal{M}|^2 \frac{\hat{s}}{(\hat{s} - M_H^2)^2 + M_H^2\Gamma_H^2}$$

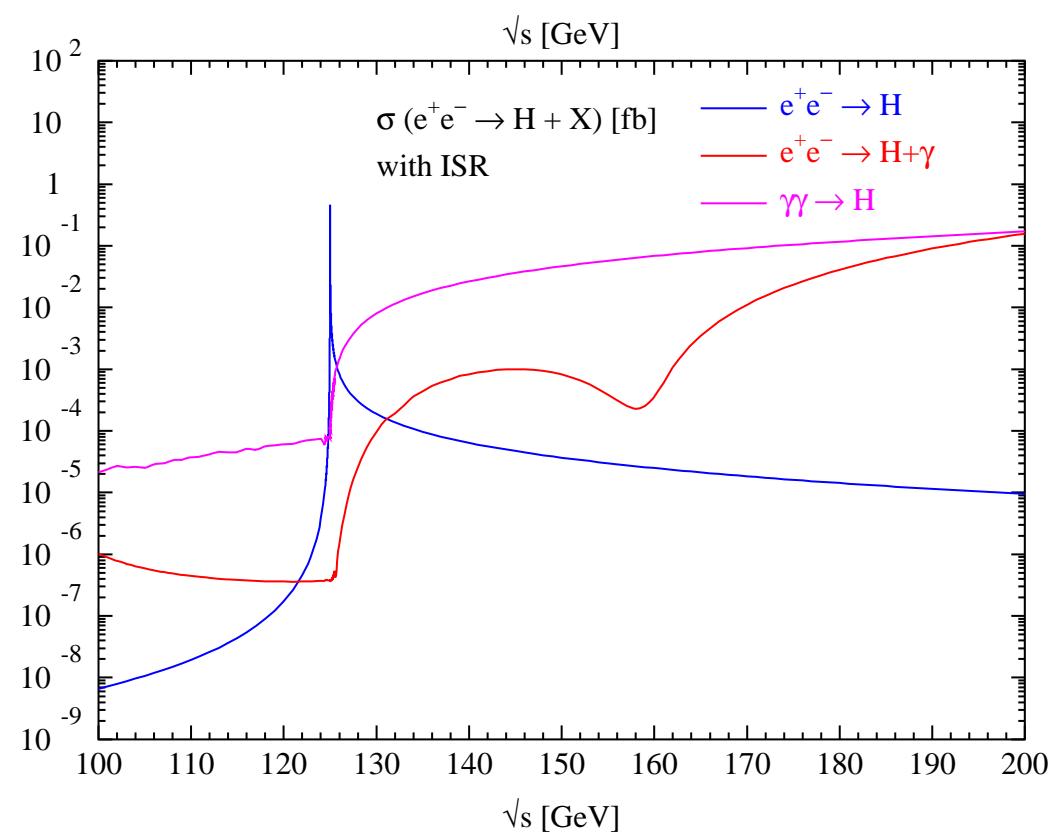
$$f_{\gamma/e}(x) = \frac{\alpha}{2\pi} \left\{ \frac{1 + (1-x)^2}{x} \log \frac{s(1-x)}{m_e^2 x^2} - 2 \frac{1-x}{x} \right\}$$

comment by E. Bagnaschi

Preliminary

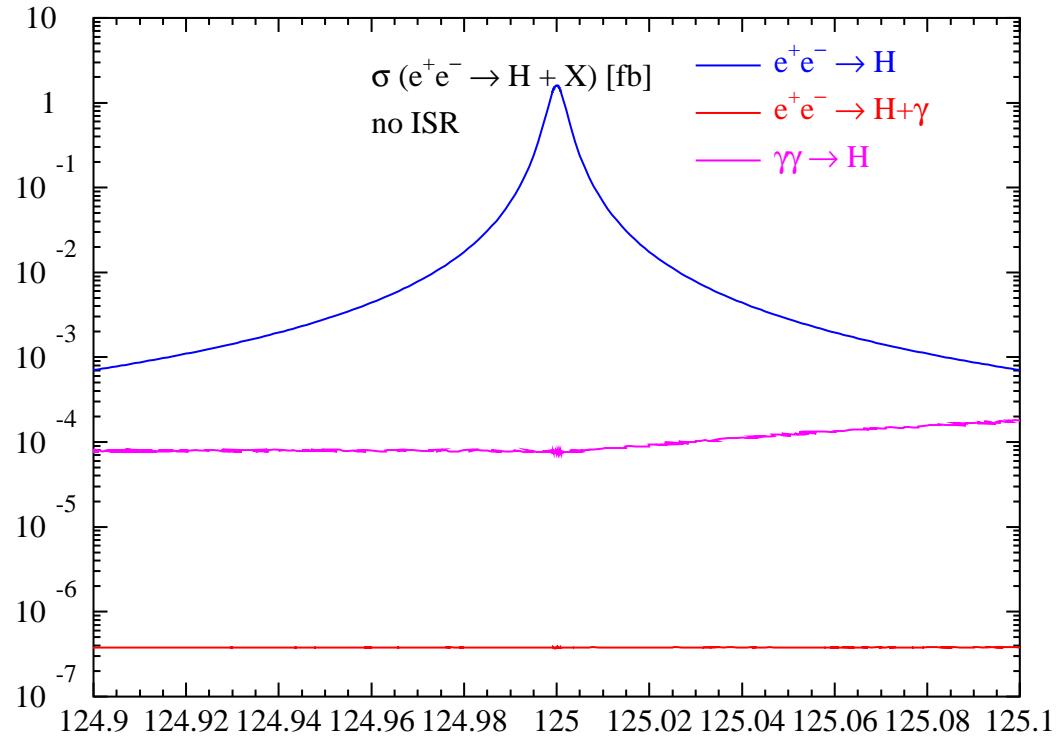


signal: no ISR

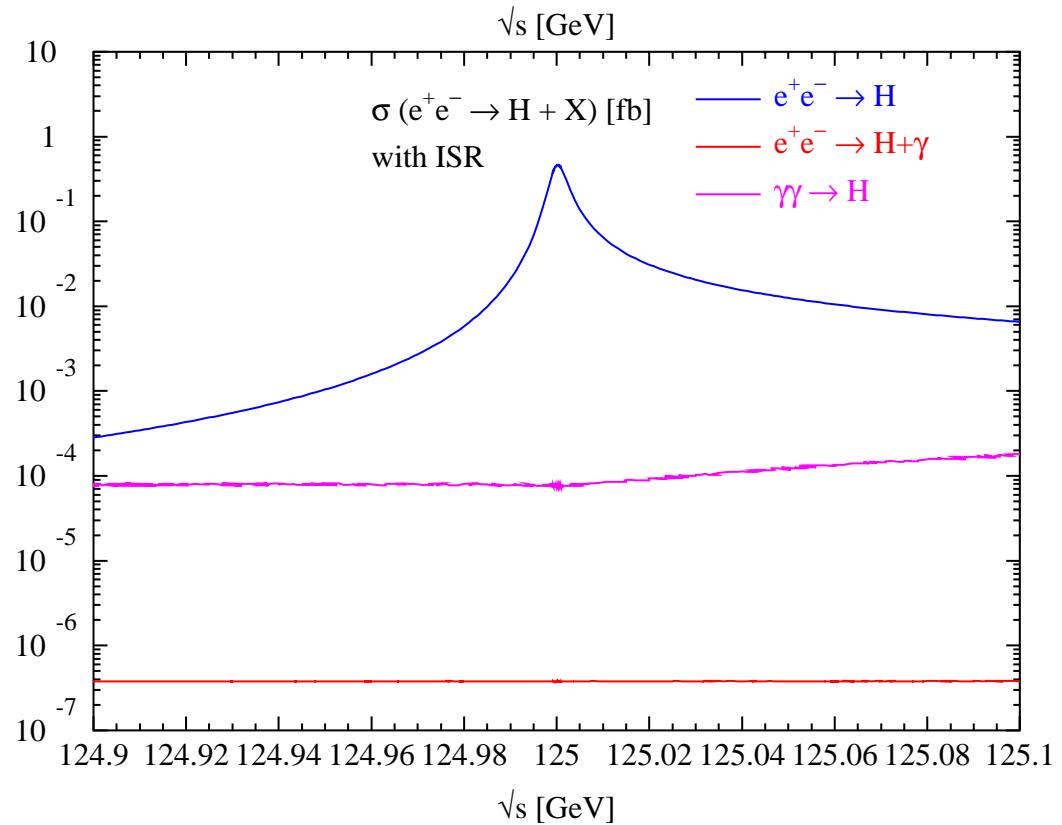


signal: with ISR

Preliminary



signal: no ISR



signal: with ISR

III *CONCLUSIONS*

- peak in $e^+e^- \rightarrow H$ large and narrow
- concurring process $e^+e^- \rightarrow H\gamma$ [t, W -loop induced] small
- concurring process $e^+e^- \rightarrow He^+e^-$ [t, W -loop induced $\gamma\gamma$ fusion] small
- huge background from non-Higgs processes \rightarrow very large luminosities required
- experimental energy spread limiting factor

BACKUP SLIDES

YFS resummation: ($\gamma_E = 0.57721566490153286060651209$)

$$\rho(\nu) = e^{\delta_{YFS}} F(\gamma) \gamma \nu^{\gamma-1} [d_s + \Delta_H(\nu)]$$

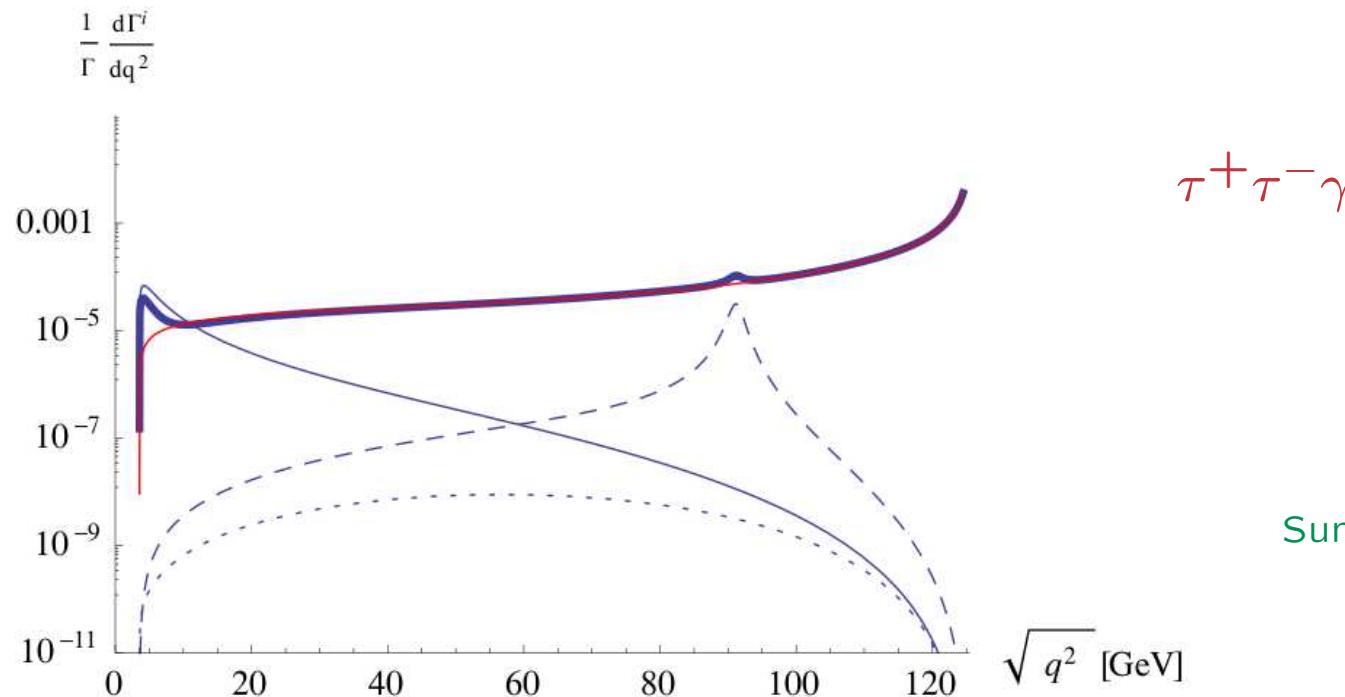
$$\gamma = 2\frac{\alpha}{\pi} \left(\log \frac{s}{m_e^2} - 1 \right)$$

$$\delta_{YFS} = \frac{\gamma}{4} + \frac{\alpha}{\pi} \left(-\frac{1}{2} + \frac{\pi^2}{3} \right)$$

$$F(\gamma) = \frac{\exp(-\gamma_E \gamma)}{\Gamma(1+\gamma)}$$

$$d_s = 1 + \frac{\gamma}{2} + \frac{\gamma^2}{8}$$

$$\Delta_H(\nu) = -\frac{\nu}{2} + \gamma \left[-\frac{\nu}{2} - \frac{1 + 3(1-\nu)^2}{8} \log(1-\nu) \right]$$



- Dalitz decays ($H \rightarrow Z\gamma \Leftrightarrow H \rightarrow \ell^+\ell^-\gamma \Leftrightarrow H \rightarrow \ell^+\ell^-$)
 [waiting for agreement ATLAS/CMS]

